

## Research Article

# Sources of Charged Higgs Pair through Double or Triple Higgs Production at Linear Colliders

Ijaz Ahmed<sup>1,2</sup>

<sup>1</sup>National Centre for Particle Physics, University of Malaya, 50603 Kuala Lumpur, Malaysia

<sup>2</sup>COMSATS Institute of Information Technology (CIIT), Islamabad 44000, Pakistan

Correspondence should be addressed to Ijaz Ahmed; [ijahmed75@gmail.com](mailto:ijahmed75@gmail.com)

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The production of triple Higgs ( $H^+H^-H^0$ ), ( $H^+H^-h^0$ ) and pairwise charged Higgs boson ( $H^+H^-$ ) is studied in the context of future linear colliders within the two-Higgs-doublet model (2HDM) type II. The aim is to compare sources of charged Higgs pair through the above processes, that is, double and triple Higgs production. Cross sections are calculated at the leading order in 2HDM type II and Minimal Supersymmetric Standard Model (MSSM). Several orders of magnitude ( $\sim 10^4$ ) enhancement are observed in 2HDM compared to MSSM, while no sizable enhancement is seen in muon collider versus electron-positron collider. The analysis is based on a heavy charged Higgs with mass above 500 GeV. It is found that double charged Higgs production cross section (being the same in 2HDM and MSSM) is few femtobarns, while the triple Higgs production cannot exceed a fraction of femtobarn within the parameter space under study.

## 1. Introduction

The main undisputed highlight of Run 1 of the Large Hadron Collider (LHC) at CERN [1, 2] is the discovery of the Standard Model (SM) Higgs boson. The measured signal strengths are quite in agreement with SM predictions. The mass of the Higgs signal is found to be near 125 GeV, which not only confirms the Higgs mechanism as a right approach towards giving masses to the electroweak particles and gauge bosons but also puts a question of possibility of existence of further Higgs bosons, as it is still not clear whether the Higgs sector is indeed minimal, containing only a single Higgs doublet. One of the straight forward ways to address such a question is simply to go beyond the SM by adding a second Higgs doublet to the field content of the model. A particularly well-motivated possibility along these lines is the Minimal Supersymmetric Standard Model (MSSM) [3] and the general (unconstrained) two-Higgs-doublet model (2HDM) [4, 5]. The general 2HDM Higgs sector contains two CP-even neutral Higgs bosons,  $h^0, H^0$ , a CP-odd (pseudo-scalar) neutral Higgs boson,  $A^0$ , and pair of charged Higgs bosons,

$H^\pm$ , whereas,  $h^0$  is the SM-like Higgs boson and is the candidate for the signal observed at LHC.

The purpose of this paper is to take into account all current constraints on the type II CP-conserving 2HDM parameter space and determine the allowed ranges of the triple and double Higgs couplings to estimate their corresponding cross sections. This work would prepare the ground for collider studies. Therefore, two types of linear colliders, that is,  $e^+e^-$  and  $\mu^+\mu^-$ , are compared for all processes. Both of these types of processes have been extensively searched during the last years at Tevatron, Large Electron-Positron Collider (LEP), and currently Large Hadron Collider (LHC). The recent results from LHC exclude a large parameter space of the light charged Higgs,  $m_{H^\pm} < 160$  GeV, if  $\text{BR}(H^\pm \rightarrow \tau\nu) = 1$  and heavy charged Higgs at  $\tan\beta > 50$  [6, 7]. Furthermore, a remarkable restriction over charged Higgs masses comes from Flavor Changing Neutral Current (FCNC), radiative B-meson decay, whose branching ratio  $\text{BR}(b \rightarrow s\gamma) \approx 3 \times 10^{-4}$  [8] is measured with sufficient precision that becomes sensitive to new physics. The charged Higgs contribution in above branching ratio increases with

decreasing  $m_{H^\pm}$ . This channel has been studied by BaBar and Belle collaborations in detail [9, 10] and the up-to-date limit excludes charged Higgs lighter than 480 GeV at 95% CL [11]. Therefore, the analysis presented throughout the paper is based on  $m_{H^\pm} \geq 500$  GeV.

A phenomenological study of triple Higgs production, including event study and the effect of neutral Higgs masses,  $m_H$  and  $m_A$ , on the cross section can be found in [12] for  $m_{H^\pm} < 500$  GeV.

In order to choose right masses for Higgs bosons of 2HDM, one should be aware that the global fit to electroweak measurements requires  $\Delta\rho$  to be  $O(10^{-3})$  [13]. This requirement does not allow large mass splitting between Higgs bosons. Therefore, we adopt degenerate masses for neutral Higgs bosons, that is,  $m_H = m_A = m_{H^\pm}$ .

Regarding the neutral Higgs mass  $m_{H^0}$ , recently CMS and ATLAS experiments have excluded a wide range of  $m_{H^0}$  in MSSM via a study of  $H \rightarrow \tau\tau$  channel at  $\sqrt{s} = 8$  TeV [14, 15]. Higgs boson decays to gauge bosons such as  $H \rightarrow WW$ ,  $ZZ$ , and  $\gamma\gamma$  have not been studied for MSSM neutral Higgs bosons. They are, however, considered for the light SM Higgs boson as different sources of Higgs boson production. Our chosen MSSM points are outside the LHC excluded area if one sets  $m_H = m_A = m_{H^\pm} \geq 500$  GeV and  $\tan\beta = 10$ .

This study is intended to be suitable at  $e^+e^-$  colliders or muon colliders. The muon collider is expected to get the integrated luminosity around  $125 \text{ fb}^{-1}$  at  $\sqrt{s} = 1.5$  TeV and  $440 \text{ fb}^{-1}$  at  $\sqrt{s} = 3$  TeV. This is a unique machine, to be designed to provide high luminosity, very small energy spread, excellent stability, and good shielding of muon beam decay backgrounds.

## 2. Theoretical Framework

The theoretical basis is a two-Higgs-doublet model with the general potential as follows [16, 17]:

$$\begin{aligned} \mathcal{V} = & m_{11}^2 \Phi_1^\dagger \Phi_1 + m_{22}^2 \Phi_2^\dagger \Phi_2 - [m_{12}^2 \Phi_1^\dagger \Phi_2 + \text{h.c.}] \\ & + \frac{1}{2} \lambda_1 (\Phi_1^\dagger \Phi_1)^2 + \frac{1}{2} \lambda_2 (\Phi_2^\dagger \Phi_2)^2 \\ & + \lambda_3 (\Phi_1^\dagger \Phi_1) (\Phi_2^\dagger \Phi_2) + \lambda_4 (\Phi_1^\dagger \Phi_2) (\Phi_2^\dagger \Phi_1) \\ & + \frac{1}{2} \lambda_5 (\Phi_1^\dagger \Phi_2)^2 \\ & + [\lambda_6 (\Phi_1^\dagger \Phi_1) + \lambda_7 (\Phi_2^\dagger \Phi_2)] (\Phi_1^\dagger \Phi_2) + \text{h.c.} \end{aligned} \quad (1)$$

The free parameters of such a model are

$$\lambda_1, \lambda_2, \lambda_3, \lambda_4, \lambda_5, \lambda_6, \lambda_7, m_{12}^2, \tan\beta \quad (2)$$

in the general basis. The CP violation or Flavor Changing Neutral Currents (FCNC) are not assumed as they are naturally suppressed via the Natural Flavor Conservation (NFC) mechanism by imposing  $Z_2$  symmetry on the Lagrangian [18, 19] leads to  $\lambda_6 = \lambda_7 = 0$ . Furthermore, we choose  $\sin(\beta - \alpha) = 1$  which takes the region of study very close to MSSM parameter space. With the above setting, that is,

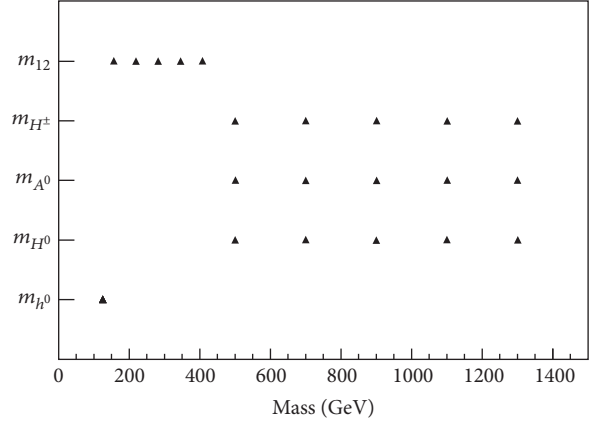


FIGURE 1: Higgs boson masses and the value of  $m_{12}$  which respect physical requirements on model potential as well as experimental limits and observations on Higgs boson masses.

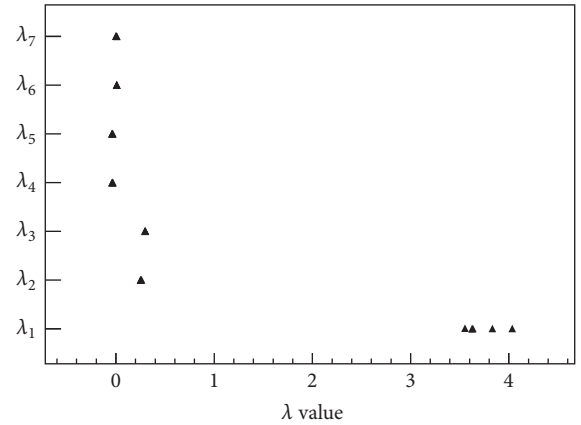


FIGURE 2: The values of  $\lambda$  parameters which result in Higgs boson masses shown in Figure 1.

$\tan\beta = 10$  and  $\sin(\beta - \alpha) = 1$ , free parameters can alternatively be taken as  $m_h$ ,  $m_H$ ,  $m_A$ ,  $m_{H^\pm}$ ,  $m_{12}^2$ ,  $\tan\beta$  because there is a correspondence between Higgs boson masses and  $\lambda$  values. We use 2HDMC package [20] to ensure that chosen parameters are consistent with current experimental limits and respect also the potential unitarity, perturbativity, and stability. A point is chosen if it satisfies all above requirements. Figures 1 and 2 show Higgs boson masses used in the analysis and their corresponding  $\lambda$  and  $m_{12}$  values extracted from 2HDMC. The choice of the default set of MSSM parameters is  $M_{\text{SUSY}} = 1000$  GeV,  $\mu = 200$  GeV,  $A_t = 1000$  GeV,  $A_b = 1000$  GeV, and  $A_\tau = 1000$  GeV, which reflect the benchmark  $m_h^{\text{max}}$  scenario.

## 3. Charged Higgs Pair Production

Double Higgs or pair production of charged Higgs bosons at linear colliders has been extensively investigated in the literature mainly in MSSM in the context of  $e^+e^-$  and  $\mu^+\mu^-$  colliders, [21–24]. The two types of charged Higgs decay have been adopted in those analyses, that is,  $H^\pm \rightarrow \tau\nu$  and

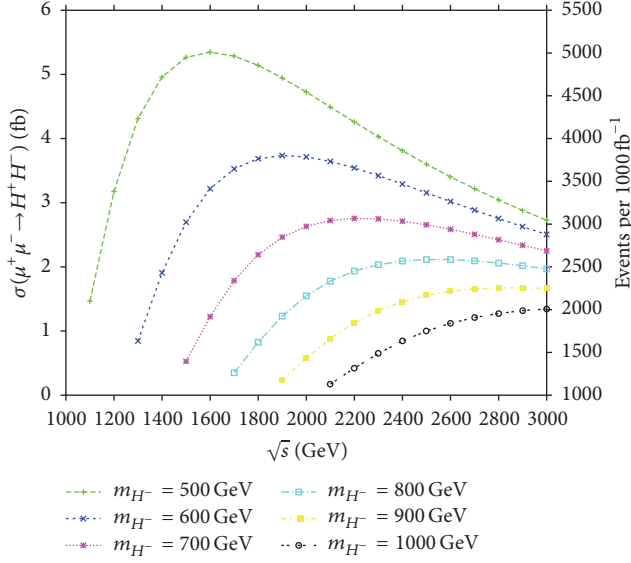


FIGURE 3: Total cross section  $\sigma(H^+H^-)$  (in fb) as a function of  $\sqrt{s}$  for the tree-level Higgs boson pair production at 2HDM or MSSM at a muon collider.

$H^\pm \rightarrow tb$ . One of the weak points of this channel is that it is limited by the center of mass energy of the collider and a charged Higgs mass heavier than  $\sqrt{s}/2$  is not produced unless a small negligible rate due to off-shell production is considered. In Figure 3, the total production cross section  $\sigma(\mu\mu \rightarrow H^+H^-)$  (in pb) is plotted as a function of center of mass energy  $\sqrt{s}$  (in GeV) for degenerate masses of Higgs bosons in 2HDM type II. The charged Higgs pair production cross section is independent of  $\tan\beta$  and  $m_A$  values. As is seen in Figure 3, the corresponding production rates reach few femtobarns, attaining several thousand events per 1000 fb<sup>1</sup> of integrated luminosity. It should be noted that there are Feynman diagrams which could result in enhancement of the production rate. They are  $s$ -channel diagrams containing neutral Higgs bosons contributions and their couplings with leptons as well as  $t$ -channel charged Higgs diagrams shown in Figure 4. However, such diagrams have very small contributions to the total cross section and the  $Z/\gamma$ -mediating diagrams dominate. Since the relevant coupling in this diagram is  $ZH^+H^-$  ( $e \cdot \cot 2\beta$ ) and  $\gamma H^+H^-$  ( $e$ ) and both of them are model independent gauge couplings, the obtained cross sections are the same in 2HDM and MSSM. Table 1 compares contribution of different diagrams in the total production cross section at a certain point in parameter space. The conclusion for this observation is that double charged Higgs production ( $H^+H^-$ ) is produced at  $\mu^+\mu^-$  and  $e^+e^-$  colliders with no sizable difference.

Since production cross section is the same at both 2HDM and MSSM, it is difficult to distinguish the two models using double Higgs ( $H^+H^-$ ) channel. It should be noted that one can look to single charged Higgs production in association with  $W$  boson. This channel has been studied in [25–27] and shows that signal cross section is larger at muon colliders and is almost independent of the charged Higgs mass and

TABLE 1: Total cross section (in fb) of double charged Higgs at  $e^+e^-$  and  $\mu^+\mu^-$  colliders at  $\sqrt{s} = 3$  TeV. The Higgs boson masses are  $m_A = m_{H^\pm} = 500$  GeV.

Collider type	$H^+H^-$					
	MSSM		2HDM			
	$Z/\gamma$	$h^0/H^0$	Total	$Z/\gamma$	$h^0/H^0$	Total
$\mu^+\mu^-$	2.73	$1.40 \times 10^{-6}$	2.73	2.74	$1.41 \times 10^{-6}$	2.74
$e^+e^-$	2.73	$1.27 \times 10^{-13}$	2.73	2.74	$7.61 \times 10^{-16}$	2.74

increases with  $\tan\beta$ . A similar analysis of this channel at LHC has been reported in [28]. In the rest of the paper, triple Higgs production is studied. It might be possible that the triple Higgs boson self-interactions play a role in this endeavor; however, as is shown in the next sections, cross section of triple Higgs production ( $H^+H^-H^0$ ,  $H^+H^-h^0$ ) is smaller than double Higgs production ( $H^+H^-$ ).

#### 4. Triple Higgs Production

The triple Higgs production has been studied in different papers in the context of linear colliders [29–33]. The production processes under consideration are  $H^+H^-H^0$  and  $H^+H^-h^0$  in  $\mu^+\mu^-$  annihilation within the 2HDM type II with their corresponding trilinear Higgs bosons couplings given in (3) and (4) together with their corresponding MSSM values in (5) and (6), where usual abbreviations such as  $s_W = \sin\theta_W$  and  $s_\beta = \sin\beta$  have been used [34]:

$$H^+H^\pm H^0 \text{ (2HDM)} : \frac{-e}{m_W \cdot s_W \cdot s_{2\beta}^2} \left( c_\beta^3 s_{2\beta} s_\alpha m_H^2 \right. \quad (3)$$

$$\left. + c_\alpha s_{2\beta} s_\beta^3 m_H^2 - 2s_{\alpha+\beta} \mu_{12}^2 + c_{\beta-\alpha} s_{2\beta}^2 m_{H^\pm}^2 \right)$$

$$H^+H^\pm h^0 \text{ (2HDM)} : \frac{e}{m_W \cdot s_W \cdot s_{2\beta}^2} \left( 2c_{\alpha+\beta} \mu_{12}^2 \right. \quad (4)$$

$$\left. - c_\alpha c_\beta^3 s_{2\beta} m_h^2 + s_{2\beta} s_\alpha s_\beta^3 m_h^2 - s_{2\beta}^2 s_{\beta-\alpha} m_{H^\pm}^2 \right)$$

$$H^+H^\pm H^0 \text{ (MSSM)} = \frac{-e \cdot m_Z^2}{2m_W s_W} \left( c_{2W} c_{\beta-\alpha} + s_{2\beta} s_{\alpha+\beta} \right) \quad (5)$$

$$H^+H^\pm h^0 \text{ (MSSM)} = \frac{-e \cdot m_Z^2}{2m_W s_W} \left( 2c_W^2 s_{\beta-\alpha} + c_{2\beta} s_{\alpha+\beta} \right). \quad (6)$$

Figure 5 shows Feynman diagrams of triple Higgs production. In this case, there is  $\tan\beta$  dependence, contrary to the case of  $H^+H^-$  production. The  $H^+H^-H^0$  coupling effectively grows as  $\tan\beta$  or  $\cot\beta$  for  $\tan\beta \gg 1$  or  $\tan\beta \ll 1$ , respectively. The corresponding cross section can vary either by  $\tan^2\beta$  at larger  $\tan\beta$  values or by  $\cot^2\beta$  at small  $\tan\beta$  values, respectively, in contrast to the situation in MSSM, where the triple Higgs coupling undergoes radiative corrections [35] and does not have any possible source of enhancement as can be seen from (5) showing tree-level coupling in MSSM. Equation (5) clearly indicates that the couplings are naturally gauge-like and hence the expected cross section remains rather small [22]. Figures 6 and 7 show cross sections of  $H^+H^-H^0$  and

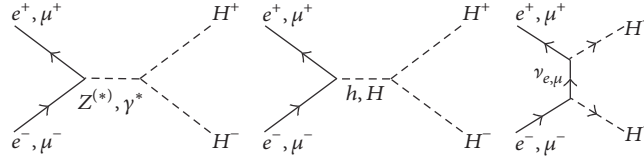


FIGURE 4: The possible Feynman diagrams for double Higgs production at lepton colliders.

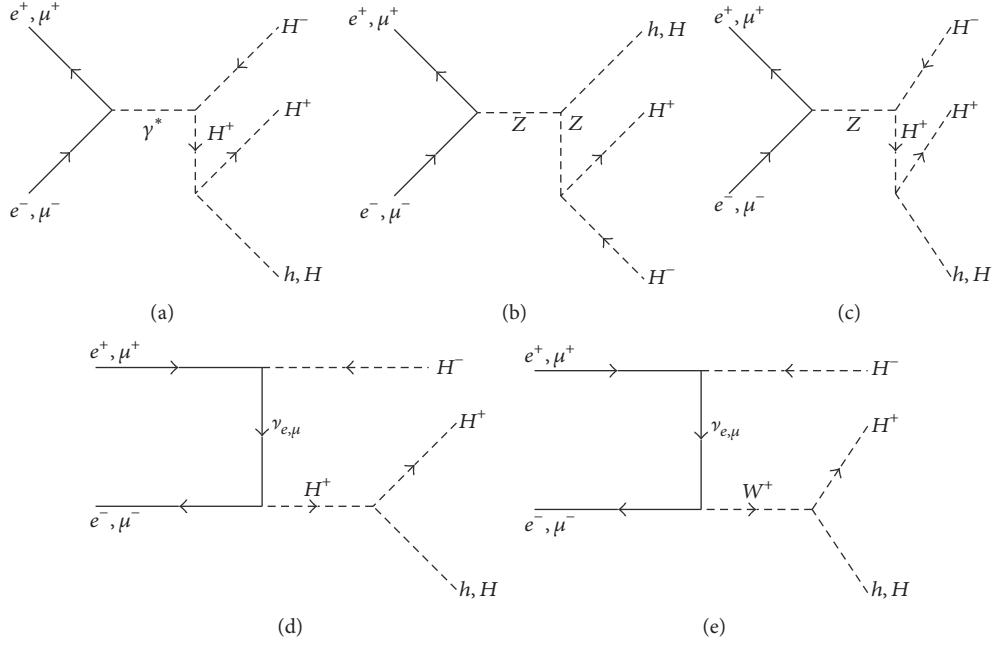
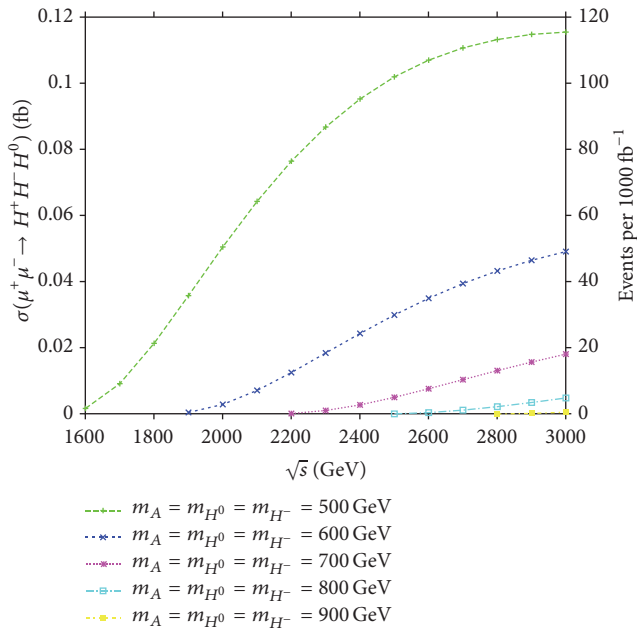
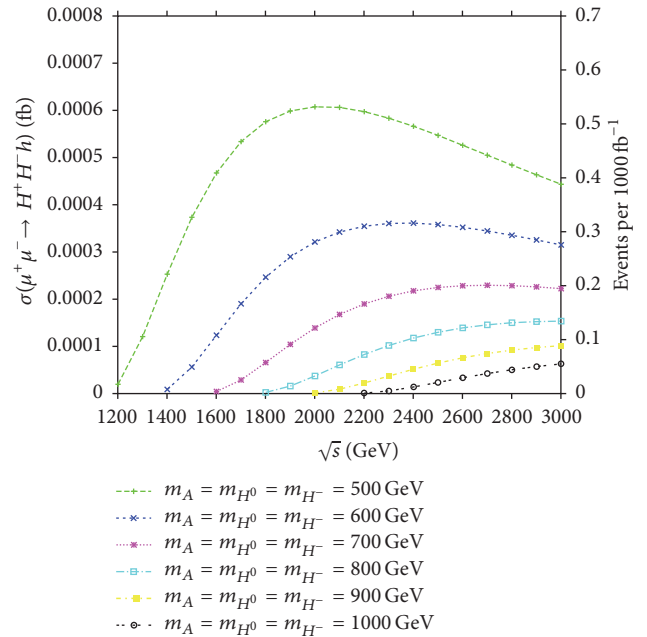


FIGURE 5: The possible Feynman diagrams for triple Higgs production at lepton colliders.

FIGURE 6: Triple Higgs production ( $H^+ H H^0$ ) within 2HDM type II for various  $m_{H^\pm}$  values. Neutral Higgs bosons have masses equal to the mass of the charged Higgs.FIGURE 7: Triple Higgs production ( $H^+ H^- h^0$ ) within 2HDM type II for various  $m_{H^\pm}$  values. Neutral Higgs bosons have masses equal to the mass of the charged Higgs.

$H^+H^-h^0$  production for different set of charged Higgs masses and center of mass energies.

One can see the contribution of different parts of the set of diagrams and compare them in a table as was done for double Higgs production. Table 2 shows such information for  $H^+H^-H^0$  and  $H^+H^-h^0$  processes, where a comprehensive comparison between 2HDM and MSSM results for  $\mu^+\mu^-$  and  $e^+e^-$  colliders has been done.

At this stage, a comparison between the two types of processes, that is, double and triple Higgs production, is performed by finding a center of mass energy which gives the maximum allowed cross section. In Table 3,  $\sigma_{\max}(H^+H^-H^0)$ ,  $\sigma_{\max}(H^+H^-h^0)$  and  $\sigma(H^+H^-)$  at two different centers of mass energies ( $\sqrt{s} = 1$  TeV and  $\sqrt{s} = 1.5$  TeV) are listed. The abbreviated name ‘‘NP’’ stands for ‘‘Not Possible,’’ due to not enough center of mass energy to produce a given event. As is seen from Table 3, the MSSM cross sections for both triple Higgs cases are extremely small reaching the largest value of  $\sigma_{\max}(H^+H^-H^0) \approx 10^{-6}$  fb and  $\sigma_{\max}(H^+H^-h^0) \approx 10^{-5}$  fb. Comparing 2HDM and MSSM cross sections, it is observed that one could reach a factor of  $10^5$  enhancement in 2HDM cross section of  $H^+H^-H^0$ , while, for  $H^+H^-h^0$ , the 2HDM cross section is typically 10 times larger than the corresponding MSSM value. On the basis of these results, the triple Higgs boson channels are generally much more promising in 2HDM than MSSM framework, but they do not reach the level of  $H^+H^-$  cross section. The double Higgs production cross section is substantially larger than triple Higgs channels either in MSSM or in 2HDM. Therefore, any study of the charged Higgs at linear colliders should be based on double Higgs production, unless a large integrated luminosity and center of mass energy are available for production of triple Higgs processes. Such events can only be considered as complementary processes to shed some light on triple Higgs couplings.

## 5. Cross Section at $e^+e^-$ and $\mu^+\mu^-$ Colliders

In this section, cross sections of three selected channels at two possible choices of colliders, that is,  $e^+e^-$  and  $\mu^+\mu^-$ , are compared. As shown in Figure 4, the charged Higgs pair production may proceed through three types of diagrams. The left diagram includes  $s$ -channel electroweak propagators, while the middle one consists of neutral Higgs bosons as propagator. The CP-odd neutral Higgs boson does not contribute in the production cross section due to having zero coupling with charged Higgs pair. Therefore, we only take into account the electroweak and CP-even neutral Higgs bosons contribution. The amount of cross section coming from each diagram may be considerably important probably after a long run of any one of these two colliders in future. There is also a  $t$ -channel process (the right diagram in Figure 4), which has negligibly small contribution at both  $e^+e^-$  and  $\mu^+\mu^-$  colliders due to the small Yukawa coupling between the charged Higgs and the leptons. Table 1 demonstrates that  $H^+H^-$  cross section is dominated by the  $Z/\gamma$ -mediating diagram (the left diagram in Figure 4) and no difference is obtained between  $e^+e^-$  and  $\mu^+\mu^-$  even by introducing

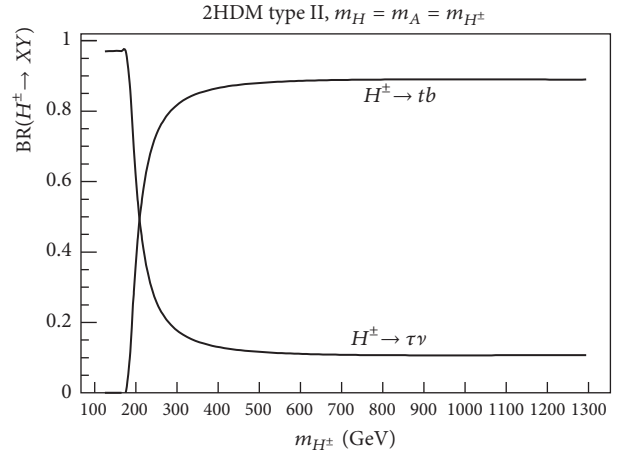


FIGURE 8: Branching ratio of charged Higgs decays in 2HDM type II.

the nonzero masses of electron and muon and their corresponding couplings in the Lagrangian. There is in fact a large enhancement in  $h/H$  mediating diagram; however, the total cross section is not affected by this enhancement due to the negligible contribution of collider dependent diagrams.

The same conclusion holds for the triple Higgs production processes shown in Figure 5. The diagram with only  $Z/\gamma$  propagator (Figure 5(b)) has the dominant contribution and, as seen in Table 2, the effect of collider dependent diagrams, that is, the ones with neutral or charged Higgs coupling with leptons, is small.

In summarizing the results, we can conclude that, for the charged Higgs searches, the most suitable channel from cross section point of view is charged Higgs pair production having largest cross section of the order of few fb in MSSM and 2HDM, while the selected triple Higgs processes ( $H^+H^-H^0$ ,  $H^+H^-h^0$ ) have very tiny cross sections. The triple Higgs searches could benefit from the neutral Higgs decay to  $b\bar{b}$  which can be identified by  $b$ -tagging algorithm. However, due to the larger particle multiplicity compared to  $H^+H^-$  production, selection efficiencies of the triple Higgs are expected to be small. The final state to analyze for both double and triple Higgs cases could be defined by selecting one of the two possible decay channels of the charged Higgs, that is,  $H^\pm \rightarrow \tau\nu$  and  $H^\pm \rightarrow tb$ . As Figure 8 shows, the charged Higgs branching ratio of decay to  $\tau\nu$  goes down to the level of 0.1, for  $m_{H^\pm} \geq 500$  GeV, while decay to  $tb$  receives a branching ratio of 0.9. Since cross sections are small, especially, for triple Higgs production, the main channel to study is expected to be the one with top and bottom quarks in the hard scattering final state. Of course in this case, a large number of particles are produced due to the decay of the top quark to  $W$  boson and its subsequent decay to leptons or light jets. One the contrary, the final state with  $\tau$  leptons involved may be identified using  $\tau$  identification algorithm, but it certainly suffers from the very tiny number of events in hand. In any case, the double Higgs production is expected to be more promising than the triple Higgs production due to the larger cross section.

TABLE 2: Total cross section (in fb) comparison between electron-positron and muon-muon colliders as well as among MSSM and 2HDM at  $\sqrt{s} = 1.5$  TeV.

Collider type	Model type	$H^+H^-H^0$			$H^+H^-h^0$		
		$Z/\gamma$	$h^0/H^0$	Total	$Z/\gamma$	$h^0/H^0$	Total
$\mu^+\mu^-$	MSSM	$1.93 \times 10^{-7}$	$3.49 \times 10^{-6}$	$3.69 \times 10^{-6}$	$1.69 \times 10^{-5}$	$2.85 \times 10^{-8}$	$1.69 \times 10^{-5}$
$\mu^+\mu^-$	2HDM	0.116	$2.24 \times 10^{-6}$	0.116	$4.44 \times 10^{-4}$	$1.90 \times 10^{-10}$	$4.44 \times 10^{-4}$
$e^+e^-$	MSSM	$1.93 \times 10^{-7}$	$7.77 \times 10^{-11}$	$1.93 \times 10^{-7}$	$1.69 \times 10^{-5}$	$6.34 \times 10^{-13}$	$1.69 \times 10^{-5}$
$e^+e^-$	2HDM	0.116	$5.10 \times 10^{-11}$	0.116	$4.44 \times 10^{-4}$	$1.03 \times 10^{-19}$	$4.44 \times 10^{-4}$

TABLE 3: The maximum cross section (in fb) for the leading order double Higgs and triple Higgs processes within MSSM and 2HDM at two different centers of mass energies  $\sqrt{s} = 1.5$  TeV and  $\sqrt{s} = 3$  TeV.

Process	$m_{H^\pm}$ GeV	$\sigma_{\max}$ (1.5 TeV)		$\sigma_{\max}$ (3 TeV)	
		MSSM	2HDM	MSSM	2HDM
$\mu^+\mu^- \rightarrow H^+H^-$	500	5.26	5.26	2.72	2.72
$\mu^+\mu^- \rightarrow H^+H^-H^0$	500	NP	NP	$3.68 \times 10^{-6}$	0.12
$\mu^+\mu^- \rightarrow H^+H^-h^0$	500	$1.44 \times 10^{-5}$	$1.69 \times 10^{-5}$	$1.69 \times 10^{-5}$	$4.44 \times 10^{-4}$
$\mu^+\mu^- \rightarrow H^+H^-$	700	5.28	5.28	2.25	2.25
$\mu^+\mu^- \rightarrow H^+H^-H^0$	700	NP	NP	$3.86 \times 10^{-7}$	$1.81 \times 10^{-2}$
$\mu^+\mu^- \rightarrow H^+H^-h^0$	700	NP	NP	$1.13 \times 10^{-5}$	$2.2 \times 10^{-4}$
$\mu^+\mu^- \rightarrow H^+H^-$	900	NP	NP	1.66	1.66
$\mu^+\mu^- \rightarrow H^+H^-H^0$	900	NP	NP	$3.86 \times 10^{-8}$	$5.77 \times 10^{-4}$
$\mu^+\mu^- \rightarrow H^+H^-h^0$	900	NP	NP	$5.27 \times 10^{-6}$	$1.02 \times 10^{-4}$

## 6. Summary

The most promising channels were studied for the charged Higgs production through double or triple Higgs production at future  $e^+e^-$  and  $\mu^+\mu^-$  linear colliders in the framework of CP-conserving two-Higgs-doublet model type II and Minimal Supersymmetric Standard Model. The sizes of double and triple Higgs production cross sections were computed both in MSSM and in 2HDM by respecting all the current experimental and phenomenological constraints on Higgs masses, requirement of perturbativity, and vacuum stability condition. The maximum cross sections at two different center of mass energies of 1.5 TeV and 3 TeV were obtained to compare these processes between MSSM and 2HDM where several orders of magnitude difference was observed. A comparison between  $e^+e^-$  and  $\mu^+\mu^-$  colliders was also performed leading to no sizable difference at both colliders. The conclusion is that double Higgs production cross section is much larger than triple Higgs, thus keeping it as the main source of the charged Higgs at linear colliders. There are diagrams which depend on the collider type (incoming lepton), but they do not have a sizable contribution to affect the total cross section. This fact is observed for both double and triple Higgs productions. The double Higgs production is less limited by the center of mass energy of the system compared to the triple Higgs production and thus can probe heavier charged Higgs bosons at a given center of mass energy. The final state to choose for a data analysis would be  $H^+H^- \rightarrow t\bar{t}b$  and  $H^+H^-H^0(H^+H^-h^0) \rightarrow t\bar{t}b\bar{b}$  for the top and bottom final states and  $H^+H^- \rightarrow \tau^+\nu\tau\bar{\nu}$  and

$H^+H^-H^0(H^+H^-h^0) \rightarrow \tau^+\nu\tau\bar{\nu}b\bar{b}$  for  $\tau$  and missing transverse energy final states. While final states with  $\tau$  and missing transverse energy are suitable for light charged Higgs searches, heavy charged Higgs with  $m_{H^\pm} \geq 500$  GeV would require a final state including top and bottom quarks.

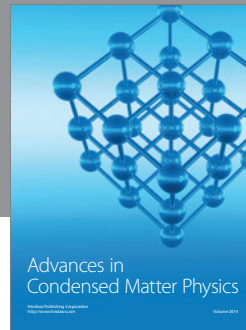
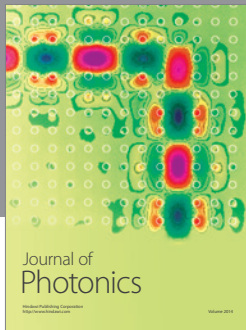
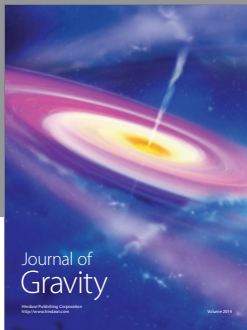
## Competing Interests

The author declares that there are no competing interests.

## References

- [1] G. Aad, T. Abajyan, B. Abbott et al., "Observation of a new particle in the search for the Standard Model Higgs boson with the ATLAS detector at the LHC," *Physics Letters B*, vol. 716, no. 1, pp. 1–29, 2012.
- [2] S. Chatrchyan, V. Khachatryan, A. M. Sirunyan et al., "Observation of a new boson at a mass of 125 GeV with the CMS experiment at the LHC," *Physics Letters B*, vol. 716, no. 1, pp. 30–61, 2012.
- [3] S. P. Martin, "A supersymmetry primer," *Advanced Series on Directions in High Energy Physics*, vol. 18, pp. 1–98, 1998.
- [4] T. D. Lee, "A theory of spontaneous  $T$  violation," *Physical Review D*, vol. 8, no. 4, pp. 1226–1239, 1973.
- [5] G. Branco, P. Ferreira, L. Lavoura, M. Rebelo, M. Sher, and J. P. Silva, "Theory and phenomenology of two-Higgs-doublet models," *Physics Reports*, vol. 516, no. 1–2, pp. 1–102, 2012.
- [6] ATLAS Collaboration, in *Proceedings of the 1st International Conference on Supersymmetry and Unification of Fundamental Interactions*, ATLAS-CONF-2013-090, Trieste, Italy, August 2013.

- [7] CMS Collaboration, “Search for charged Higgs bosons with the  $H^+$  to tau nu decay channel in the fully hadronic final state at  $\sqrt{s} = 8$  TeV,” CMS PAS HIG-14-020, 2014.
- [8] W. M. Yao et al., “Review of particle physics,” *Journal of Physics G: Nuclear and Particle Physics*, vol. 33, no. 1, pp. 1–1232, 2006.
- [9] J. P. Lees, V. Poireau, V. Tisserand et al., “Measurement of  $\mathcal{B}(B \rightarrow X_s \gamma)$  photon energy spectrum, and the direct  $CP$  asymmetry in  $B \rightarrow X_{s+d} \gamma$ ,” *Physical Review D*, vol. 86, no. 11, Article ID 112008, 2012.
- [10] A. Limosani, H. Aihara, K. Arinstein et al., “Measurement of inclusive radiative  $B$ -meson decays with a photon energy threshold of 1.7 GeV,” *Physical Review Letters*, vol. 103, no. 24, Article ID 241801, 2009.
- [11] M. Misiak, H. M. Asatrian, R. Boughezal et al., “Updated next-to-next-to-leading-order QCD predictions for the weak radiative  $B$ -meson decays,” *Physical Review Letters*, vol. 114, no. 22, Article ID 221801, 2015.
- [12] M. Hashemi and I. Ahmed, “Observability of triple or double charged Higgs production in two Higgs doublet model type II at an  $e^+e^-$  linear collider,” *International Journal of Modern Physics A*, vol. 30, Article ID 1550022, 2015.
- [13] J. Beringer, J.-F. Arguin, R. M. Barnett et al., “Review of particle physics,” *Physical Review D*, vol. 86, no. 1, Article ID 010001, 2012.
- [14] CMS Collaboration, “Search for neutral MSSM Higgs bosons decaying to a pair of tau leptons in  $pp$  collisions,” *Journal of High Energy Physics*, vol. 2014, no. 10, article 160, 2014.
- [15] G. Aad, B. Abbott, J. Abdallah et al., “Search for neutral Higgs bosons of the minimal supersymmetric standard model in  $pp$  collisions at  $\sqrt{s} = 8$  TeV with the ATLAS detector,” *Journal of High Energy Physics*, vol. 2014, article 56, 2014.
- [16] S. Davidson and H. E. Haber, “Basis-independent methods for the two-Higgs-doublet model,” *Physical Review D*, vol. 72, Article ID 035004, 2005.
- [17] H. E. Haber and D. O’Neil, “Basis-independent methods for the two-Higgs-doublet model. II. The significance of  $\tan\beta$ ,” *Physical Review D*, vol. 74, no. 1, Article ID 015018, 27 pages, 2006.
- [18] S. L. Glashow and S. Weinberg, “Natural conservation laws for neutral currents,” *Physical Review D*, vol. 15, no. 7, pp. 1958–1965, 1977.
- [19] I. F. Ginzburg and M. Krawczyk, “Symmetries of two Higgs doublet model and  $CP$  violation,” *Physical Review D*, vol. 72, no. 11, Article ID 115013, 26 pages, 2005.
- [20] D. Eriksson, J. Rathsmann, and O. Stål, “2HDMC—two-Higgs-doublet model calculator,” *Computer Physics Communications*, vol. 181, no. 1, pp. 189–205, 2010.
- [21] A. Djouadi, J. Kalinowski, and P. M. Zerwas, “Exploring the SUSY Higgs sector at  $e^+e^-$  linear colliders: a synopsis,” *Zeitschrift für Physik C Particles and Fields*, vol. 57, no. 4, pp. 569–583, 1993.
- [22] A. Djouadi, H. Haber, and P. Zerwas, “Multiple production of neutral Higgs bosons at high-energy  $e^+e^-$  colliders,” *Physics Letters B*, vol. 375, no. 1–4, pp. 203–212, 1996.
- [23] S. M. Hashemi and I. S. Nefedov, “Wideband perfect absorption in arrays of tilted carbon nanotubes,” *Physical Review B*, vol. 86, no. 19, Article ID 195411, 2012.
- [24] M. Hashemi, “Charged higgs pair production in a general two higgs doublet model at  $e^+e^-$  and  $\mu^+\mu^-$  linear colliders,” *Communications in Theoretical Physics*, vol. 61, no. 1, pp. 69–74, 2014.
- [25] A. G. Akeroyd, A. Arhrib, and C. Dove, “Charged and pseudoscalar Higgs boson production at a muon collider,” *Physical Review D*, vol. 61, no. 7, Article ID 071702, 2000.
- [26] A. Akeroyd and S. Baek, “Single charged Higgs production as a probe of  $CP$  violation at a muon collider,” *Physics Letters B*, vol. 500, no. 1–2, pp. 142–148, 2001.
- [27] M. Hashemi, “Charged Higgs observability through associated production with  $W^\pm$  at a muon collider,” *The European Physical Journal C*, vol. 72, no. 11, article no. 2207, 2012.
- [28] M. Hashemi, “Possibility of observing an MSSM charged Higgs boson in association with a  $W$  boson at the LHC,” *Physical Review D*, vol. 83, no. 5, Article ID 055004, 12 pages, 2011.
- [29] S. Kanemura, Y. Okada, and E. Senaha, “Electroweak baryogenesis and quantum corrections to the triple Higgs boson coupling,” *Physics Letters B*, vol. 606, no. 3–4, pp. 361–366, 2005.
- [30] A. Arhrib, R. Benbrik, and C. Chiang, “Probing triple Higgs couplings of the two Higgs doublet model at a linear collider,” *Physical Review D*, vol. 77, no. 11, 2008.
- [31] G. Ferrera, J. Guasch, D. López-Val, and J. Solà, “Triple Higgs boson production in the linear collider,” *Physics Letters B*, vol. 659, no. 1–2, pp. 297–307, 2008.
- [32] D. López-Val and J. Solà, “Neutral Higgs-pair production at linear colliders within the general two-Higgs-doublet model: quantum effects and triple Higgs boson self-interactions,” *Physical Review D*, vol. 81, no. 3, Article ID 033003, 2010.
- [33] A. Gutiérrez-Rodríguez, A. Hernández-Ruiz, O. Sampayo, A. Chubykalo, and A. Espinoza-Garrido, “The triple Higgs Boson self-coupling at future linear  $e^+e^-$  Colliders Energies: ILC and CLIC,” *Journal of the Physical Society of Japan*, vol. 77, no. 9, Article ID 094101, 8 pages, 2008.
- [34] M. N. Dubinin and A. Semenov, “Triple and quartic interactions of higgs bosons in the general two-higgs-doublet model,” <https://arxiv.org/abs/hep-ph/9812246>.
- [35] W. Hollik and S. Peñaranda, “Yukawa coupling quantum corrections to the self-couplings of the lightest MSSM Higgs boson,” *The European Physical Journal C*, vol. 23, no. 1, pp. 163–172, 2002.



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